

Neutralino dark matter in MSSM with CPV

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LAPTH-Annecy

PLAN

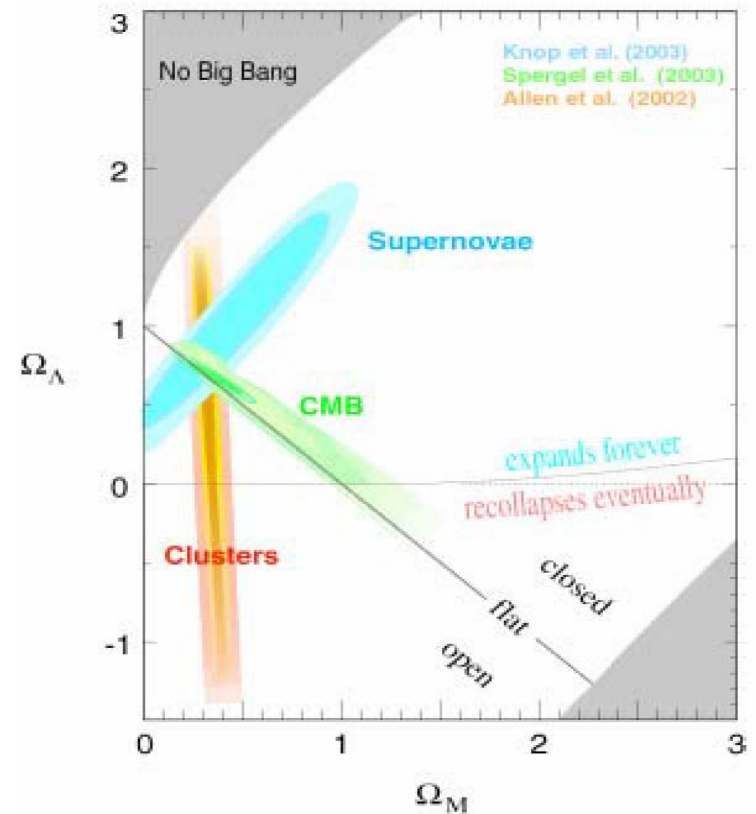
- Relic density of DM
- The CPV-MSSM
- Dark matter in the CPV-MSSM : impact of phases
- Some remarks on indirect detection
- Conclusion

Based on

GB, Boudjema, Kraml, Pukhov, Semenov, Phys.Rev.D60 (2006)

Dark matter: measurements

- In recent years : new precise determination of cosmological parameters
- Data from CMB (WMAP) agree with the one from clusters and supernovae
 - Dark matter: $23 \pm 4\%$
 - Baryons: $4 \pm 0.4\%$
 - Dark energy $73 \pm 4\%$
 - Neutrinos $< 1\%$

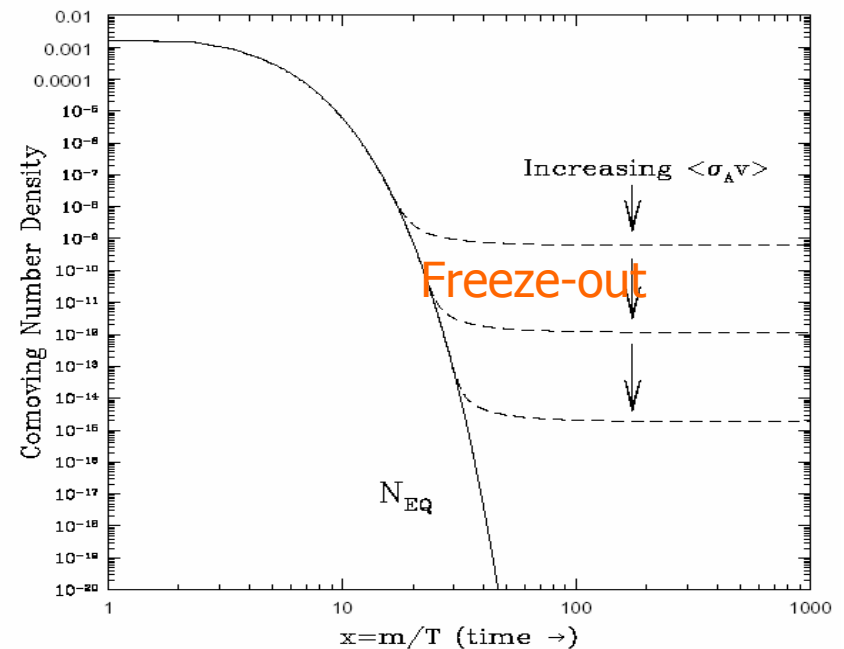


Relic density of SUSY DM

- The precise measurement of WMAP $0.095 < \Omega h^2 < .128$ provides strong constraints on candidates for DM in particular neutralino.
- At the moment predictions for Ωh^2 in MSSM varies by orders of magnitude (unknown parameters)
- Colliders will search for SUSY, might discover and measure properties of sparticles
 - refine predictions for relic density
 - see whether fits with measurements or is it necessary to modify cosmological model.

Relic density of wimps

- In early universe WIMPs are present in large number and they are in thermal equilibrium
- As the universe expanded and cooled their density is reduced through pair annihilation
- Eventually density is too low for annihilation process to keep up with expansion rate
 - Freeze-out temperature
- LSP decouples from standard model particles, density depends only on expansion rate of the universe



$$\frac{dn}{dt} = -3Hn - \langle \sigma v \rangle [n^2 - n_{eq}^2]$$

Coannihilation

- If $M(\text{NLSP}) \sim M(\text{LSP})$ then $\chi + X \rightarrow \chi' + Y$ maintains thermal equilibrium between NLSP-LSP even after SUSY particles decouple from standard ones
- Relic density depends on rate for all processes involving LSP/NLSP \rightarrow SM
- All particles eventually decay into LSP, calculation of relic density requires summing over all possible processes

$$\langle \sigma v \rangle = \frac{\sum_{i,j} g_i g_j \int_{(m_i+m_j)^2} ds \sqrt{s} K_1(\sqrt{s}/T) p_{ij}^2 \sigma_{ij}(s)}{2T \left(\sum_i g_i m_i^2 K_2(m_i/T) \right)^2}$$

Exp(- ΔM)/T

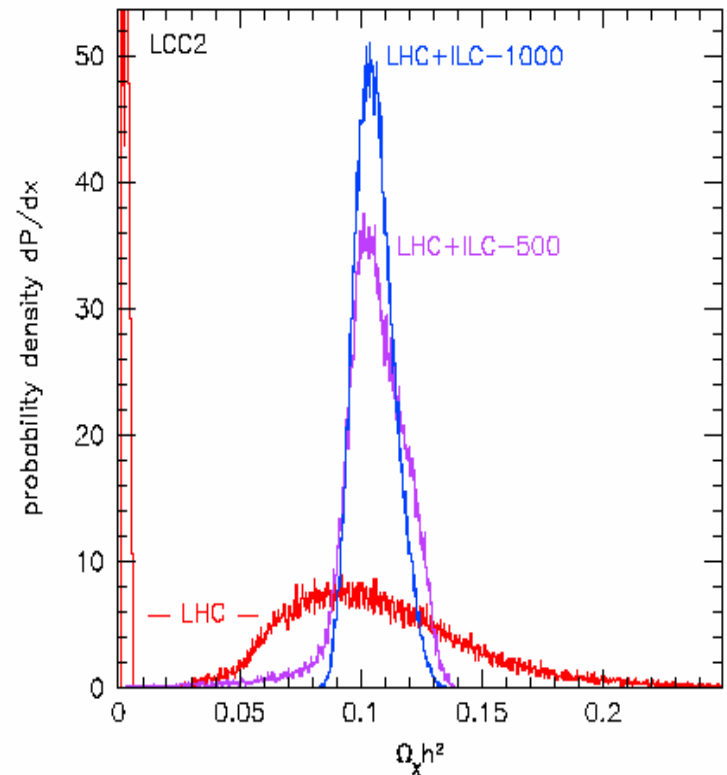
- Important processes are those involving particles close in mass to LSP
- Public codes to calculate relic density: micrOMEGAs, DarkSUSY, IsaRED

Probing cosmology using collider information

- Whether prediction for relic density can have sufficient precision (within the context of a given model) to be useful (at the level of WMAP(15%) or PLANCK (6%)) depends on the SUSY or New physics scenario.
- Many parameters enter computation of relic density, only a handful of relevant ones for each scenario – work is going on in North America, Asia and Europe to determine both the LHC and ILC potential
 - Moroi, Bambade, Richard, Zhang, Martyn, Tovey, Polesello, Lari, D. Zerwas, Allanach, Belanger, Boudjema, Pukhov, Battaglia, Birkedal, Gray, Matchev, Alexander, Fields, Hertz, Jones, Meyraiban, Pivarski, Peskin, Dutta, Kamon, Arnowitt, Khotilovith, Nojiri...

An example in MSSM ... bino/Higgsino LSP

- Improving precision on prediction of relic density using collider information
- Recent study of determination of parameters and reconstruction of relic density in a few scenarios: e.g. bino/Higgsino
- LHC: not enough precision
- ILC: chargino pair production sensitive to bino/Higgsino mixing parameter
- ILC: roughly 15% precision on Ωh^2



Baltz et al hep-ph/0602187

Precise prediction of relic density...

- But this assumes that all SUSY parameters are real...
- Phases can induce large shifts in relic density.

CPV-MSSM

- Phases are generic
- MSSM with phases might explain baryonic asymmetry of the universe-electroweak baryogenesis
 - Need strong first order phase transition
 - Finite temperature effective potential
 - $V = AT^2\phi^2 - ET\phi^3 + \lambda\phi^4$, condition : $2E/\lambda > 1$
 - In SM requires Higgs mass < 50 GeV
 - New physics solution: light RH stops in MSSM + light neutralino/chargino + CP (Carena et al)

CPV-MSSM –The model

- Parameters of the model (universality)

$M_1, \mu, \tan\beta, M_{H^\pm}, M_S, A_t, A_\tau = 1\text{TeV}$

- Phases

- $\Phi(\mu)$: strongly constrained by edm
- $\Phi(M_1)$
- $\Phi(M_3)$: mostly irrelevant for relic
- $\Phi(A_t)$: Higgs mass and mixing + also stop sector (if stop NLSP)
- $\Phi(A_\tau)$: relevant for relic if stau NLSP
- $\Phi(A_e)$: irrelevant for relic, important for edm

Neutralino LSP

- Prediction for relic density depend on parameters of model

- Mass of neutralino LSP
- Nature of neutralino** : determine the coupling to Z, h, A ...

- $M_1 < M_2 < \mu$ bino
- $\mu < M_1, M_2$ Higgsino
- $M_2 < M_1 < \mu$ Wino

$$\begin{pmatrix} M_1 & 0 & -m_Z \cos \beta \sin \theta_W & m_Z \sin \beta \sin \theta_W \\ 0 & M_2 & m_Z \cos \beta \cos \theta_W & -m_Z \sin \beta \cos \theta_W \\ -m_Z \cos \beta \sin \theta_W & m_Z \cos \beta \cos \theta_W & 0 & -\mu \\ m_Z \sin \beta \sin \theta_W & -m_Z \sin \beta \cos \theta_W & -\mu & 0 \end{pmatrix}$$

$$\tilde{\chi}_1^0 = Z_{11} \tilde{B}^0 + Z_{12} \tilde{W}_3^0 + Z_{13} \tilde{H}_d^0 + Z_{14} \tilde{H}_u^0 .$$

CPV-MSSM

- Higgs sector: mixing between 3 states (h,H,A), no longer eigenstates of CP: h1,h2,h3
- Neutralino couplings to Higgs

$$\mathcal{L}_{\tilde{\chi}_1^0 \tilde{\chi}_1^0 h_i} = -\frac{g}{2} \sum_{i=1}^3 \overline{\tilde{\chi}_1^0} (g_{h_i \tilde{\chi}_1^0 \tilde{\chi}_1^0}^S + i\gamma_5 g_{h_i \tilde{\chi}_1^0 \tilde{\chi}_1^0}^P) \tilde{\chi}_1^0 h_i$$

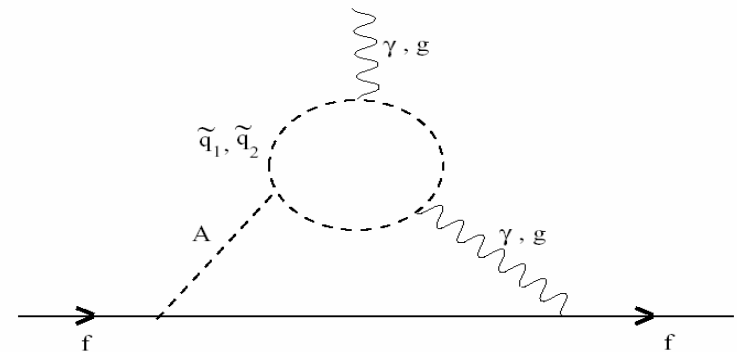
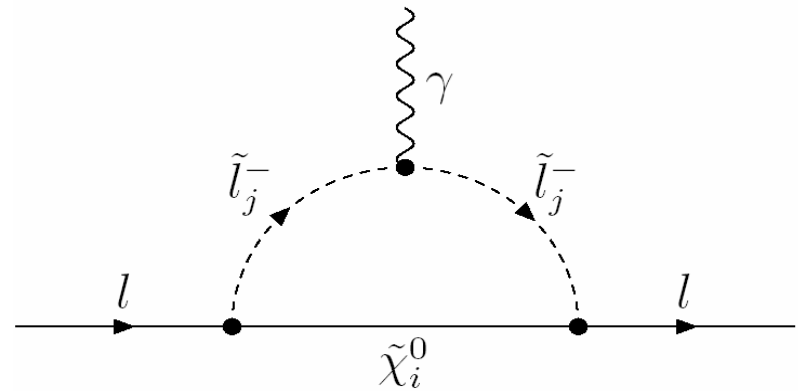
$$g_{h_i \tilde{\chi}_1^0 \tilde{\chi}_1^0}^S = \text{Re}[(N_{12}^* - t_W N_{11}^*)(H_{1i} N_{13}^* - H_{2i} N_{14}^* - iH_{3i}(s_\beta N_{13}^* - c_\beta N_{14}^*))]$$

- Neutralino-Chargino W couplings

$$\mathcal{L}_{\tilde{\chi}_j^0 \tilde{\chi}_i^+ W^-} = g W_\mu^- \overline{\tilde{\chi}_j^0} \gamma^\mu (O_{ji}^L P_L + O_{ji}^R P_R) \tilde{\chi}_i^+ + \text{h.c.} \quad O_{ji}^L = N_{j2} V_{i1}^* - \frac{1}{\sqrt{2}} N_{j4} V_{i2}^*$$

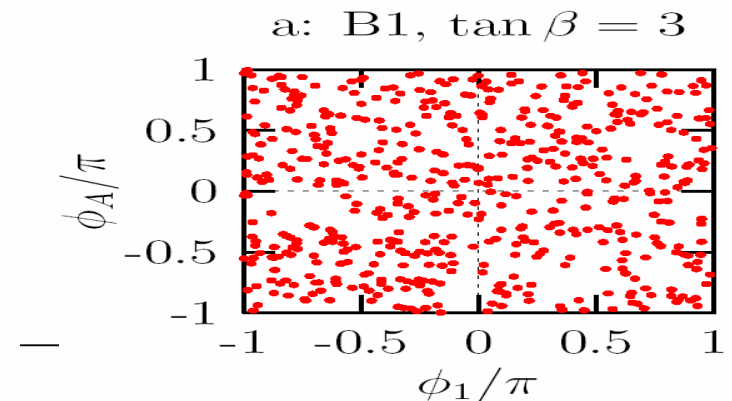
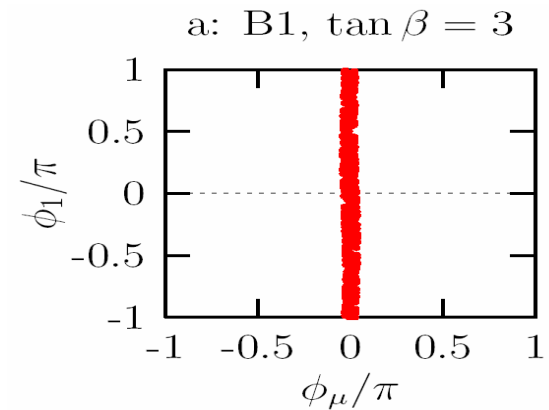
Electric dipole moment

- One loop
chargino/neutralino
- Two-loops
t,b,stop,sbottom
 - Can be dominant at large $\tan\beta$
- At small $\tan\beta$, strong constraints on Φ_μ , much less constraints on Φ_1 Φ_t



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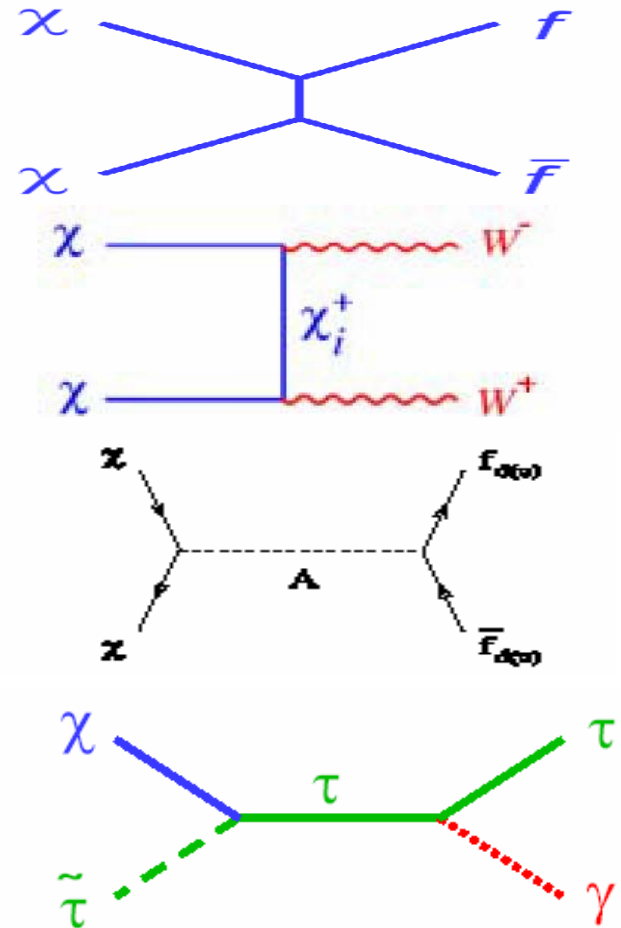
Choi, Drees, Gassmaier (2004)

Neutralino dark matter - MSSM

- bino LSP
- Bino/Higgsino
- Higgs resonance
- Sfermion coannihilation

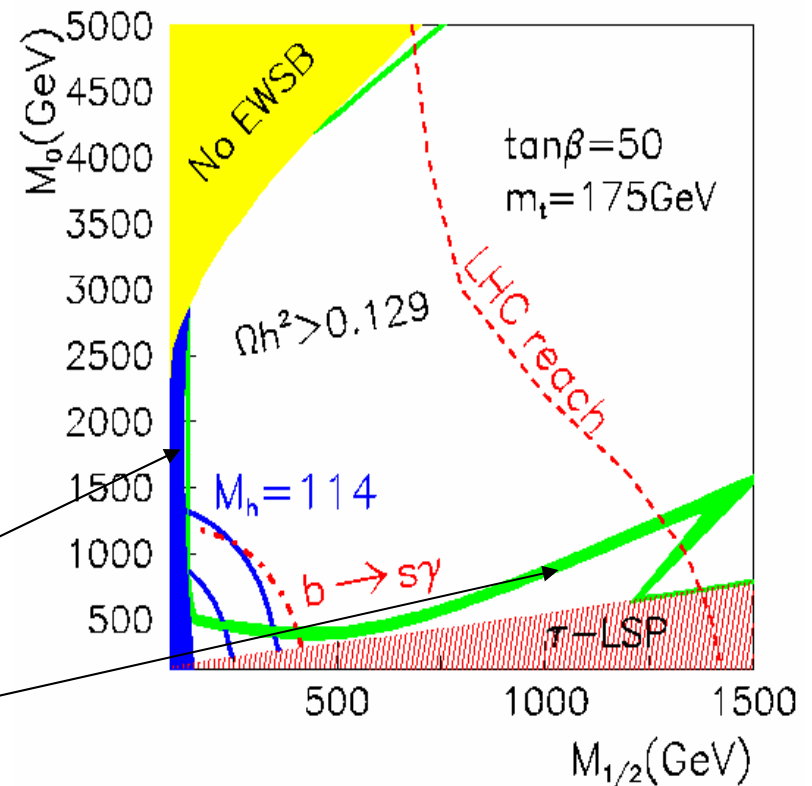
Neutralino annihilation

- Bino LSP annihilation into $f\bar{f}$
 - $\sigma \sim m_\chi^2/m_f^4$
- Mixed bino-Higgsino (wino)
 - Coupling depends on Z_{12}, Z_{13}, Z_{14} , mixing of LSP
- Annihilation near resonance (Higgs)
 - Need some mixing with Higgsino
- Co-annihilation
 - Bino LSP (sfermion coannihilation)
 - Higgsino LSP- coannihilation with chargino and neutralinos



Neutralino in mSUGRA- WMAP

- **bino – LSP**
 - In most of mSUGRA parameter space
 - Works well for light sparticles but hard to reconcile with LEP/Higgs limit
- **Mixed Bino-Higgsino**
 - Annihilation into W pairs+coannihilation
 - More likely at large $\tan\beta$
- **Resonance (Z, light/heavy Higgs)**
 - LEP constraints for light Higgs/Z
 - Heavy Higgs at large $\tan\beta$ (enhanced Hbb vertex)
- **Sfermion coannihilation**
 - Staus or stops
 - More efficient, can go to higher masses



Dark matter in CPV-MSSM

- Few studies of relic density of dark matter in MSSM with CP violation
 - Falk Olive Srednicki (1995), Nihei Sasagawa (2004), Argyrou et al (2004), Gomez et al (2005), Gondolo Freese (2002) Choi Kim (2006)
- Some report several orders of magnitude shifts in relic density due to phases -> these effects are mostly due to kinematic, e.g. LSP mass closer to Higgs resonance
- After taking into account the effects from shifts in masses, impact of phases on relic are more modest (still much larger than the uncertainty from WMAP).

CPV-MSSM & micrOMEGAs2.0

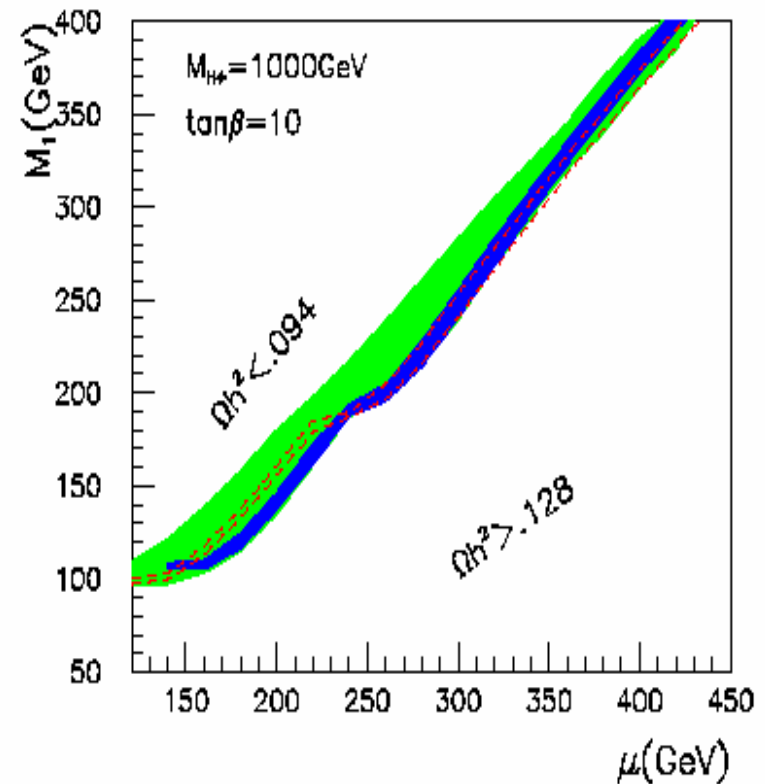
- Modify the model file in CalCHEP
- Complete tree-level matrix elements for all subprocesses from CalCHEP
- Automatically Include all possible annihilation and coannihilation channels
- Calculates the relic density for any LSP
- Solution of evolution equation and calculation of relic density with non-relativistic thermal averaging and proper treatment of poles and thresholds (Gondolo, Gelmini, NPB 360 (1991)145)
- Automatically check for presence of resonances and improves the accuracy near pole
- Interface to CPSuperH (J.S. Lee et al) : masses of Higgs and third generation and parameters of effective Higgs potential
- edm constraints are included

Neutralino dark matter – CPV-MSSM

- Same mechanisms for (co)-annihilation
- bino LSP
 - Phase itself does not impact much Ωh^2 (10-20%)
 - Higgs mass constraint is relaxed
 - If light sfermions exist LHC should tell us
- bino/Higgsino
- Higgs resonance
- Sfermion coannihilation

Bino/Higgsino

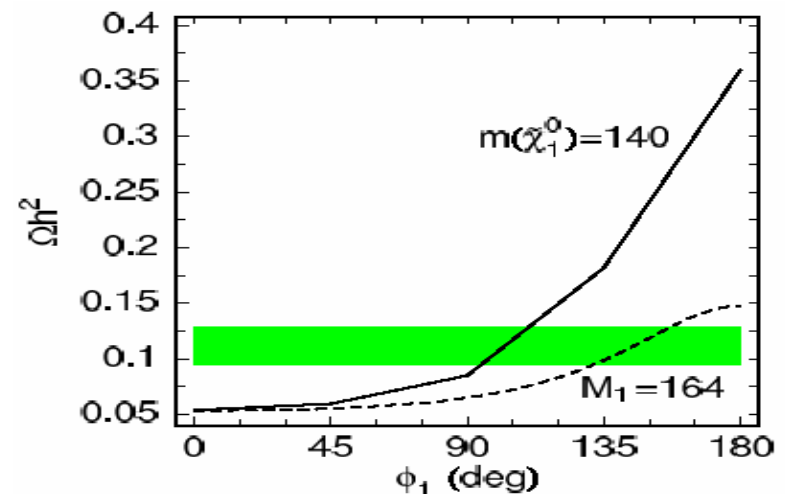
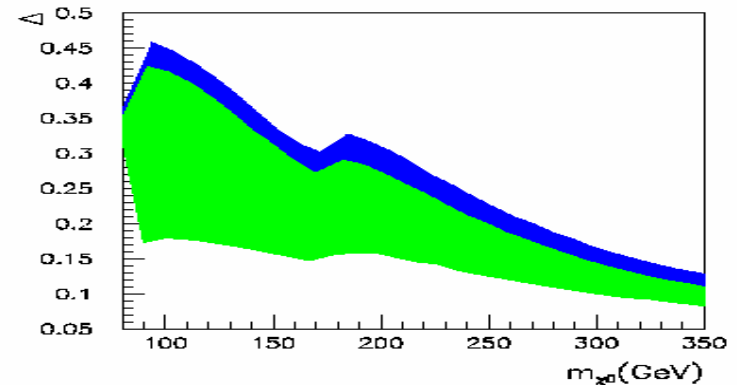
- With arbitrary phases, new allowed region where Ωh^2 below WMAP (a)
- $M(\text{LSP})$ increases with φ_1
 - more chargino coannihilation
- $\chi^0 \chi^+ W$ coupling decreases with $\varphi_1 \rightarrow$ smaller cross-section
 - Larger Higgsino fraction allowed



$\mu = 500 \text{ GeV}, M_S = 1 \text{ TeV}$

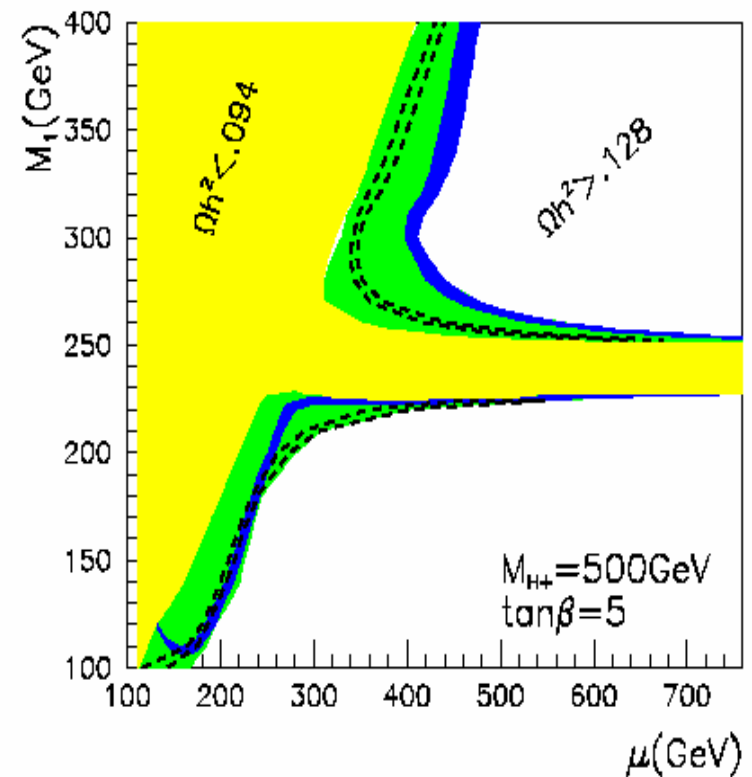
Bino/Higgsino

- Smaller mass splitting than expected in MSSM
- In this scenario once mass effects are taken into account \rightarrow Shift in Ωh^2 more important



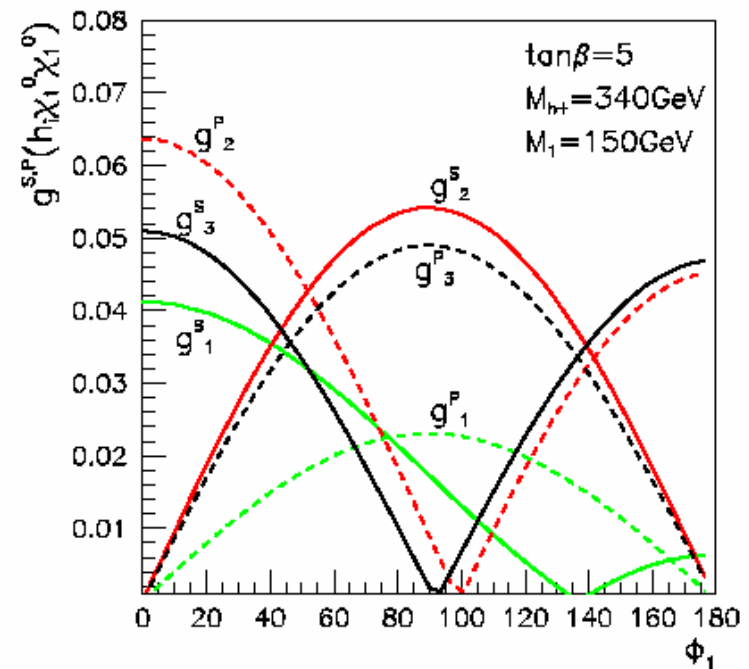
Lower $M_{H^\pm}=500\text{GeV}$

- Higgs funnel (large μ OK)
- Allowed region varying all phases
- In bino/Higgsino much wider WMAP allowed region especially near Higgs resonance
 - much larger Higgsino fraction is allowed
 - Sensitivity to mass difference $2m_\chi - m_h$



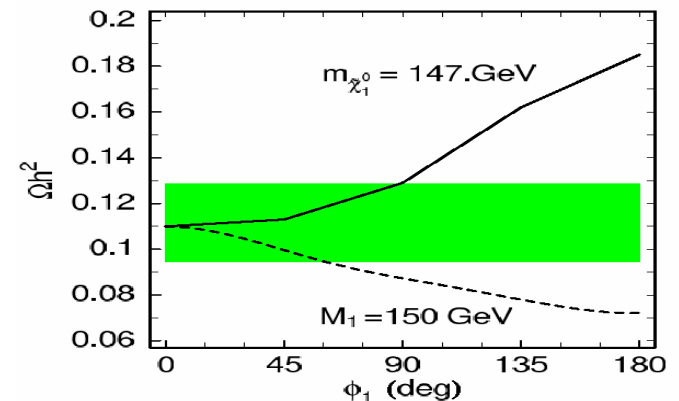
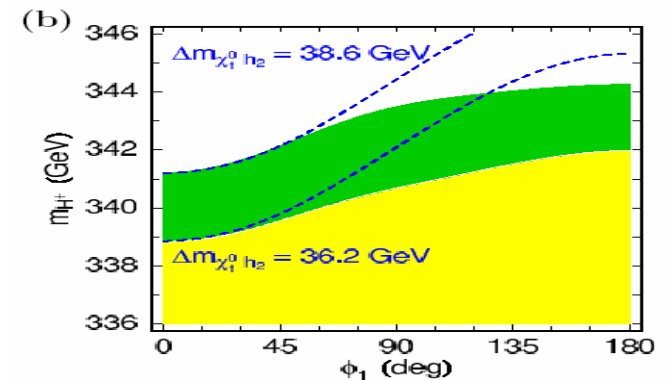
Annihilation near Higgs

- Neutralino are Majorana, at $v \rightarrow 0$ only pseudoscalar exchange, thermal averaging some scalar contribution
- Phase affects $\chi\chi h$ couplings and relative contribution of h_2 and h_3
- Phase affects mass difference $2m_\chi - m_h$



Annihilation near Higgs

- Contours of constant Ωh^2 do not follow precisely contours of constant mass difference
- Once readjusting the LSP mass so that $2m_{\tilde{\chi}} - m_h$ is constant- large corrections with ϕ_1
- Also large shifts in masses with ϕ_t , even when keeping mass difference constant shifts in Ωh^2 can be one order of magnitude.

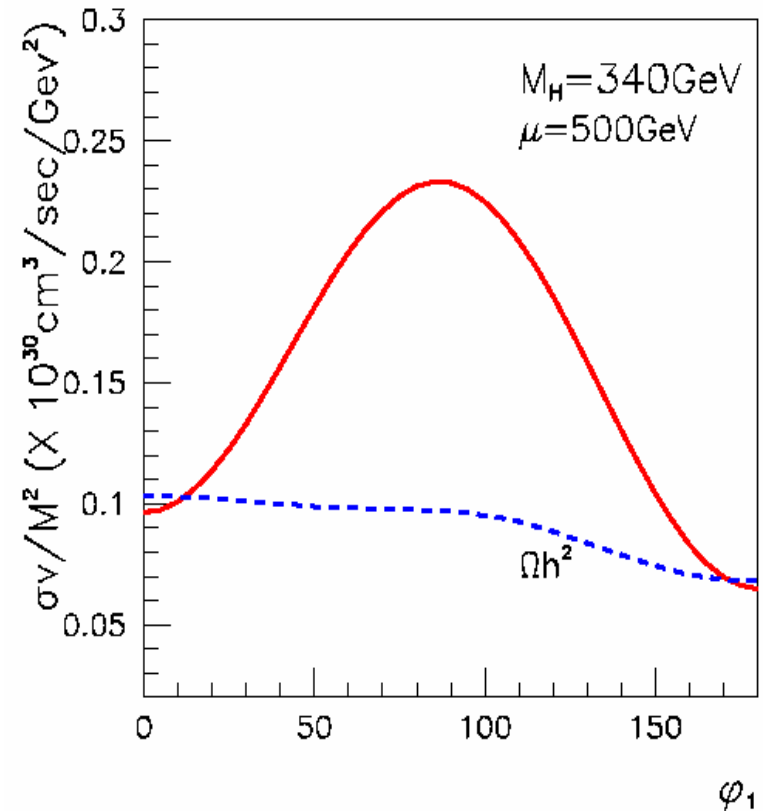


Indirect detection

- Phases should also have effect on rate for indirect detection (annihilation of neutralinos through $\gamma, e^+, \text{anti-}p$)
- Depends on σv , with $v \sim 0.001$
- But $\Omega \sim 1/\langle \sigma v \rangle$
- When $v \rightarrow 0$ only pseudoscalar contribution for Majorana neutralinos
- Effect of phase expected to be larger than for Ω
 - Especially when phase changes pseudoscalar content of h

Indirect detection

- Expect when Ω increases, σv decreases
- Often the case but not always
 - Example: Annihilation near Higgs resonance
 - Interference effect $\chi\chi$ -hh



GB, Boudjema, Kraml, Pukhov in progress

Conclusions

- Phases have strong effect on relic density up to one order of magnitude
- To have a precise prediction of DM relic density important to know precisely the masses but also measure precisely couplings and in particular phases.
- How well can this be done in scenarios interesting for DM remains to be seen
- Phases also impact indirect detection (and direct detection)