

Mirage Mediation

(Mixed Modulus - Anomaly Mediation)

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KC , A. Falkowski , H. P. Nilles , M. Olechowski

NPB 718 , 113 (2005)

KC , K. S. Jeong , K. I. Okumura

JHEP 0509 , 039 (2005)

KC , K. S. Jeong , T. Kobayashi , K. I. Okumura

PLB 633 , 355 (2006)

& also in preparation.

- Motivation

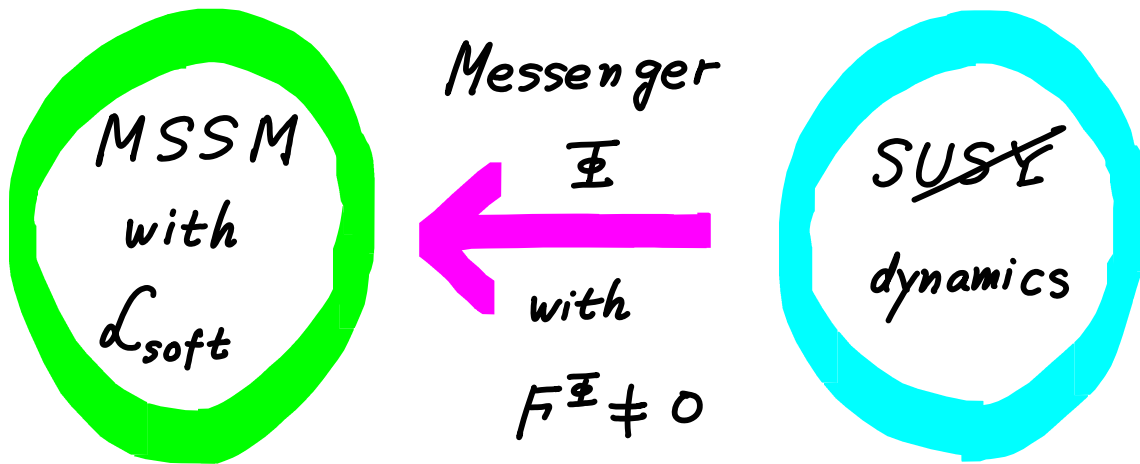
Suggested (or predicted) by large class of string compactifications stabilizing moduli at phenomenologically viable dS vacuum.

- Characteristic Features

- ★ Highly distinctive low energy spectrum of superparticles : compressed spectrum
- ★ Naturally avoid dangerous CP & $FCNC$ & also ameliorate (even avoid) the little fine-tuning for EWSB
- ★ Various possibilities for LSP DM
(Next talk by K. Okumura)

◆ Soft ~~SUSY~~ terms in models with low energy SUSY

$$-\frac{1}{2} M_a \lambda^a \lambda^a - \frac{1}{2} m_i^2 |\phi^i|^2 - \frac{1}{6} A_{ijk} y_{ijk} \phi^i \phi^j \phi^k + h.c$$



The pattern of superparticle masses is closely connected to the mechanism of messenger stabilization which determines $\langle \Xi \rangle$ & $\langle F\Xi \rangle$.

◆ Plausible messengers of ~~SUSY~~
in compactified string theory

- 4D SUGRA multiplet $\{g_{\mu\nu}, M, \dots\}$

Chiral compensator $C = 1 + M\theta^2$

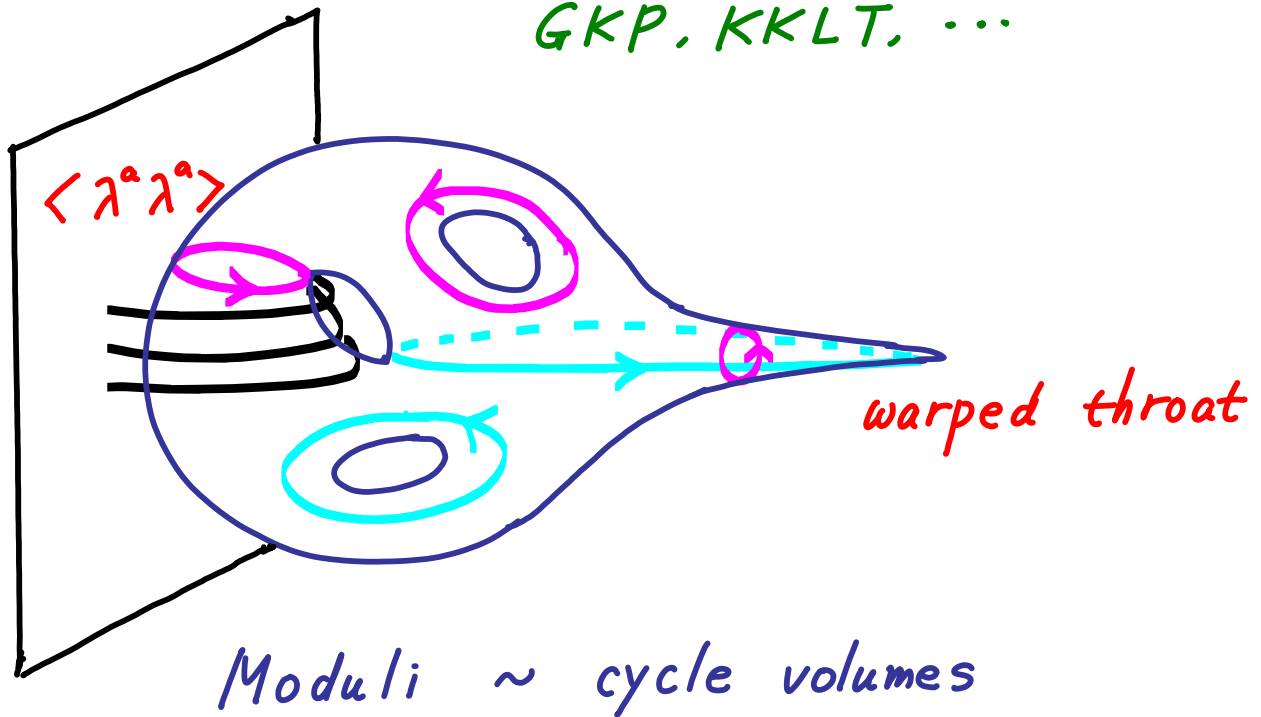
→ Anomaly mediation

- Moduli / Dilaton multiplet

→ For a reliable computation of m_{soft} in 4D EFT of string theory, one needs to know how the dilaton/moduli are stabilized at a vacuum with nearly vanishing cosmological constant.

◆ Moduli stabilization by flux & NP effects

GKP, KKLT, ...



- cycles with flux :

$$\text{flux density} \propto 1/\text{Volume}$$

- cycles with wrapping D-branes :

$$\text{gaugino condensation} \propto e^{-1/\text{Volume}}$$

(or brane instanton)

→ Fix the cycle volumes ~ moduli

- Generically some fluxes generate *warped throat* & the warp factor at the end of throat can be exponentially small :

$$e^A = \exp \left[- \frac{2\pi \int H_3}{g_{st} \int F_3} \right]$$

- Typically moduli are stabilized at a *SUSY AdS vacuum*.

$$D_{\bar{z}} W(\bar{z}_0) = 0 \quad , \quad W(\bar{z}_0) \neq 0$$

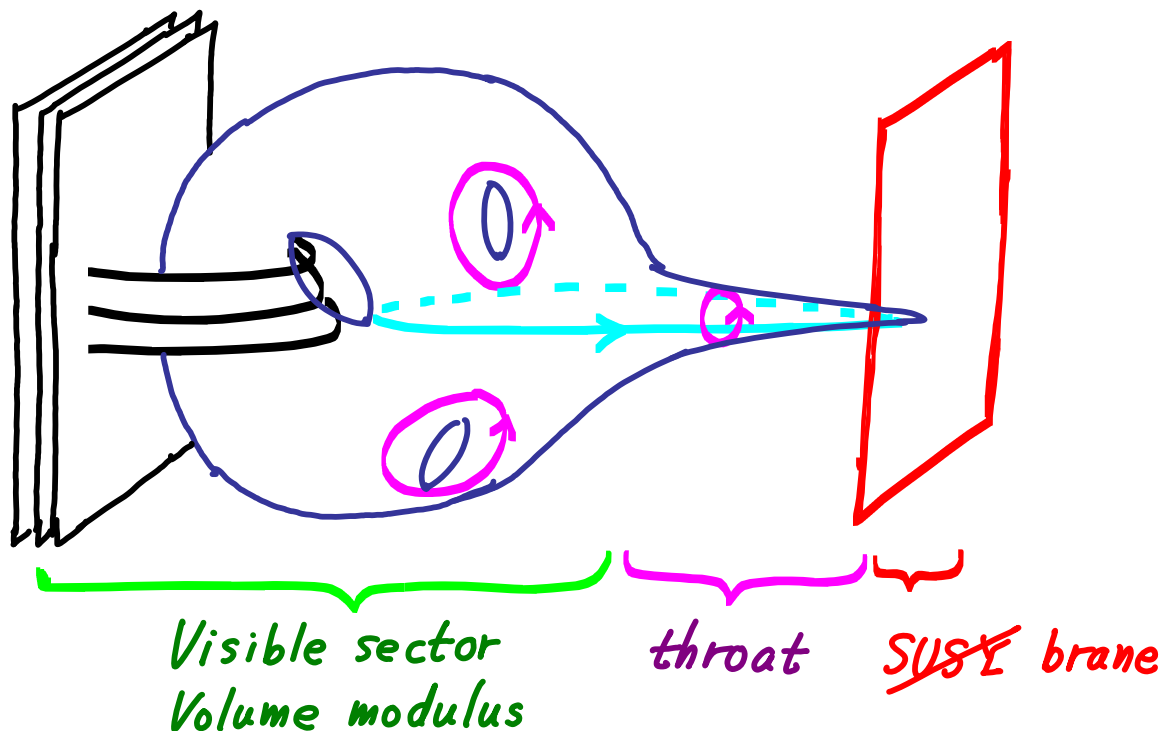
$D_{\bar{z}} W = 0$ is always a stationary point of $V_{SUGRA} = e^K [K^{\bar{z}z} |D_{\bar{z}} W|^2 - 3 |W|^2]$.

$$\rightarrow \langle V_{SUGRA} \rangle = -3 e^K |W|^2 = -3 m_{3/2}^2 M_{Pl}^2 < 0$$

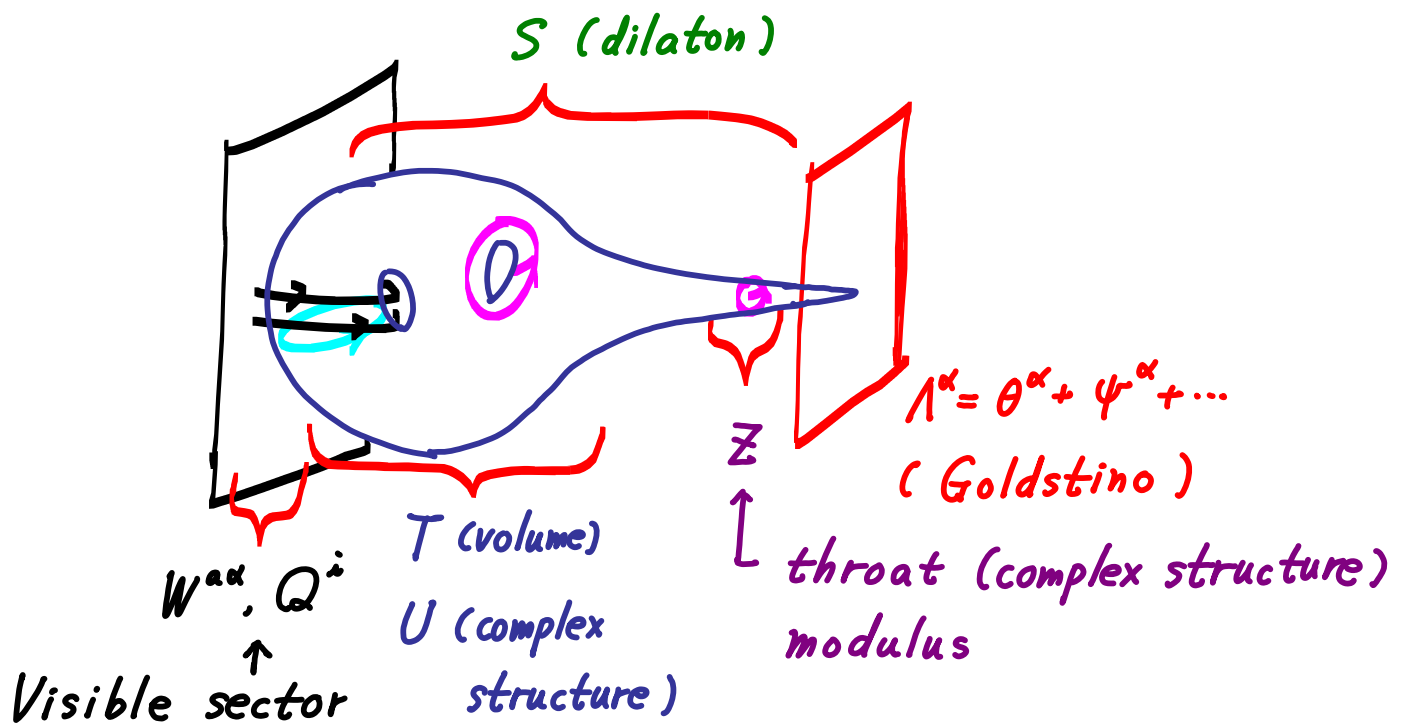
- We need to add ~~SUSY~~ brane with positive energy density.

$\overline{D3}$ (KKLT) or any brane with ~~SUSY~~ dynamics

- It is energetically favored that ~~SUSY~~ brane is stabilized at the end of warped throat.
- On the other hand, visible sector branes should be stabilized at unwarped region to achieved $M_{GUT} \sim 2 \times 10^{16} \text{ GeV}$.



◆ Example *KC, Jeong, Kobayashi, Okumura*



$$W = W_{flux} + W_{np}$$

$$\begin{aligned} W_{flux} &= \int (F_3 - 2\pi i S H_3) \wedge \Omega \\ &= \frac{n}{2\pi i} Z \ln Z - 2\pi i m S Z + \overline{\mathcal{F}}(U) - 2\pi i S \mathcal{H}(U) \end{aligned}$$

$$\begin{aligned} W_{np} &= \sum_I A_I e^{-8\pi^2 (k_I T + l_I S)} \quad (A_I \sim 1) \\ &= A_1 e^{-8\pi^2 S} + A_2 e^{-8\pi^2 (k T - l S)} \end{aligned}$$

(k, l, m, n are rational numbers)

Consider $\mathcal{F}$ & $\mathcal{H}$ with common zero :

$$\mathcal{F}(U_0) = \mathcal{H}(U_0) = 0$$

($\mathcal{F}$ & $\mathcal{H}$ are polynomials with quantized coefficients.)

S, U & Z are stabilized by W_{flux} at

$$\langle Z \rangle \simeq Z_0 \equiv e^{-\left(\frac{4\pi^2 m S_0}{n} + 1\right)}$$

$$\langle U \rangle \simeq U_0, \quad \langle S \rangle \simeq S_0 \equiv \frac{\mathcal{F}'(U_0)}{2\pi n \mathcal{H}'(U_0)}$$

with $m_{S,U,Z} \gg m_T, m_{3/2}$.

To get low energy SUSY ($m_{3/2} \sim \text{TeV}$)

together with vanishing C.C, we need to

choose the discrete parameters as

$$m_{3/2}/M_{Pl} \sim e^{-8\pi^2 S_0} \sim 10^{-14}$$

$$Z_0 \sim e^{-4\pi^2 m S_0/n} \sim \left(m_{3/2}/M_{Pl} \right)^{3/2} \sim 10^{-21}$$

$\downarrow \quad m=3n$

$\uparrow (\text{red-shift factor at the end of throat})^3$

Heavy S, U, Z can be integrate out, yielding an effective SUGRA of $T, \Lambda^\alpha, W^{\alpha\alpha}, Q^i$:

$$\mathcal{L}_{\text{eff}} = \int d^4\theta \, C\bar{C} \left[-3e^{-K/3} + \underbrace{C\bar{C} Z_0^{4/3} \theta^2 \bar{\theta}^2}_{\text{red-shift by warping}} \right] + \int d^2\theta \left[\frac{1}{4} f_a W^{\alpha\alpha} W_\alpha^\alpha + C^3 W_{\text{eff}} \right] \quad \begin{array}{l} \text{uplifting potential} \\ \text{from SUSY brane} \end{array}$$

$$K = K_0(T, \bar{T}) + Z_i(T, \bar{T}) Q^i \bar{Q}^i$$

$$f_a = T \quad \begin{array}{l} Z_0^{1/3} \sim e^{-4\pi^2 S_0} \sim 10^{-7} \\ m_{3/2} \sim e^{-8\pi^2 S_0} \sim 10^{-14} \\ (\text{in the unit with } M_{\text{Pl}} = 1) \end{array}$$

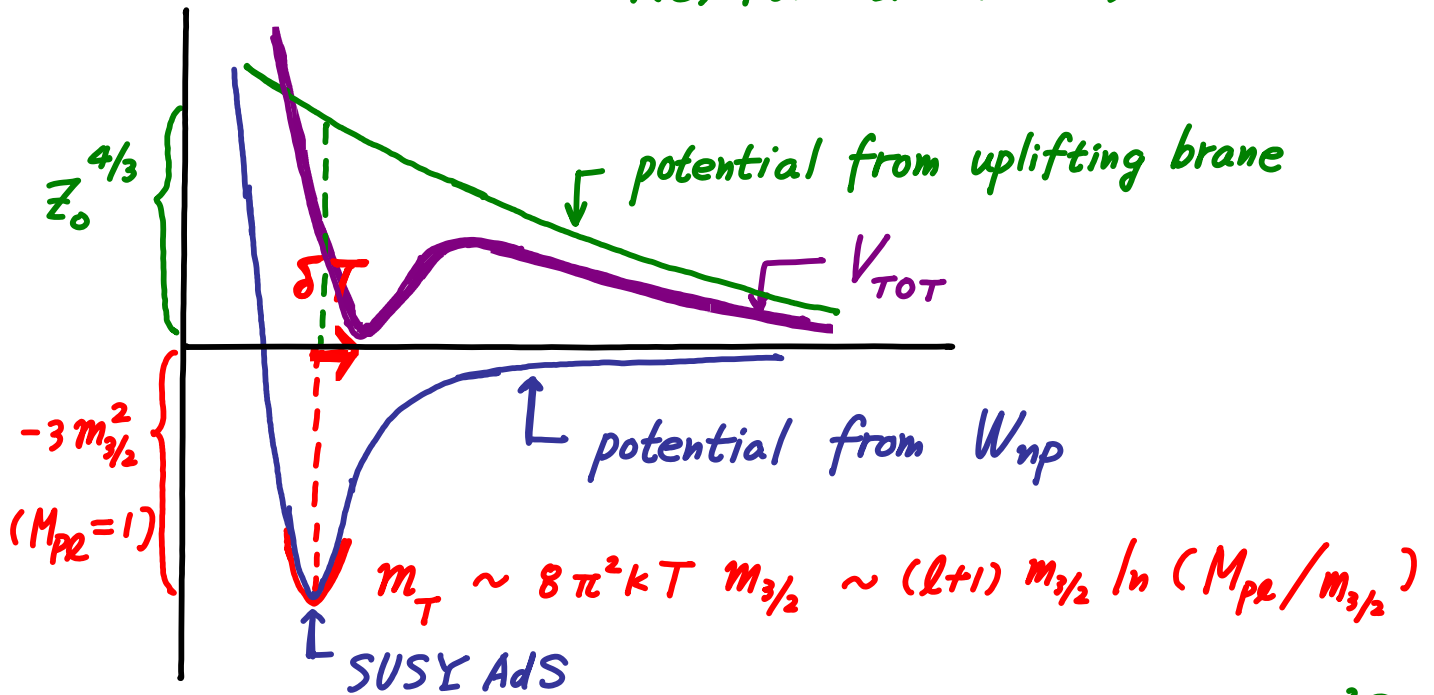
$$\begin{aligned} W_{\text{eff}} &= e^{-8\pi^2 S_0} + e^{-8\pi^2(kT - \ell S_0)} + \frac{1}{6} \lambda_{ijk} Q^i Q^j Q^k \\ &= m_{3/2} + m_{3/2}^{-\ell} e^{-8\pi^2 k T} + \frac{1}{6} \lambda_{ijk} Q^i Q^j Q^k \end{aligned}$$

The non-perturbative superpotential stabilizing the lightest modulus is directly connected to the weak to Planck scale hierarchy $m_{3/2}/M_{\text{Pl}}$:

$$e^{-8\pi^2 k T} \sim \left(m_{3/2}/M_{\text{Pl}} \right)^{\ell+1}$$

SUSY by uplifting potential

KC, Falkowski, Nilles, Olechowski



For nearly vanishing C.C ,

$$\frac{m_{3/2}}{M_{Pl}} \sim e^{-8\pi^2 S_0}$$

$$Z_0^{4/3} \sim e^{-16\pi^2 S_0 m_{3/2}} \sim \left(\frac{m_{3/2}}{M_{Pl}} \right)^2$$

SUSY shift of T : $\delta T \sim \frac{m_{3/2}^2}{m_T^2}$

modulus mediated gaugino mass

$$\rightarrow \frac{F T}{T + \bar{T}} \equiv M_0 \sim \frac{m_{3/2}^2}{m_T} \sim \frac{m_{3/2}}{(l+1) \ln(M_{Pl}/m_{3/2})}$$

Big hierarchy

$$(M_{Pl}/m_{3/2}) \sim e^{4\pi^2}$$

Little hierarchy

$$M_0 \sim \frac{m_{3/2}}{8\pi^2} \sim \frac{m_T}{(8\pi^2)^2}$$

◆ Mirage mediation pattern of soft masses

KC, Jeong, Okumura

Stabilization of the lightest set of moduli
by W_{np} & uplifting of vacuum by red-shifted
~~SUSY~~ brane

→ Moduli mediation $\sim$ Anomaly mediation

$$\uparrow \frac{F^T}{T + \bar{T}} \equiv M_0$$

$$\uparrow \frac{m_{3/2}}{8\pi^2}$$

$$M_a (M_{GUT}^{(-)}) = M_0 + \frac{b_a}{8\pi^2} g_{GUT}^2 m_{3/2}$$

$$A_{ijk} (M_{GUT}^{(-)}) = a M_0 - \frac{1}{16\pi^2} (\gamma_i + \gamma_j + \gamma_k) m_{3/2}$$

$$m_i^2 (M_{GUT}^{(-)}) = \underbrace{c M_0^2}_{\text{moduli mediation}} + \underbrace{\frac{1}{8\pi^2} \frac{\partial \gamma_i}{\partial \ln T} M_0 m_{3/2}}_{\text{mixed mediation}} - \underbrace{\frac{\gamma_i}{16\pi^2} m_{3/2}^2}_{\text{anomaly mediation}}$$

Low energy superparticle masses

$$M_a(\mu) = M_0 \left[1 - \frac{b_a}{4\pi^2} g_a^2(\mu) \ln \left(\frac{M_{\text{mirage}}}{\mu} \right) \right]$$

$$m_i^2(\mu) = M_0^2 \left[c + \frac{1}{4\pi^2} \left(\delta_i - \frac{1}{2} \dot{\delta}_i \right) \ln \left(\frac{M_{\text{mirage}}}{\mu} \right) \right]$$

$\uparrow$ 1st & 2nd generations

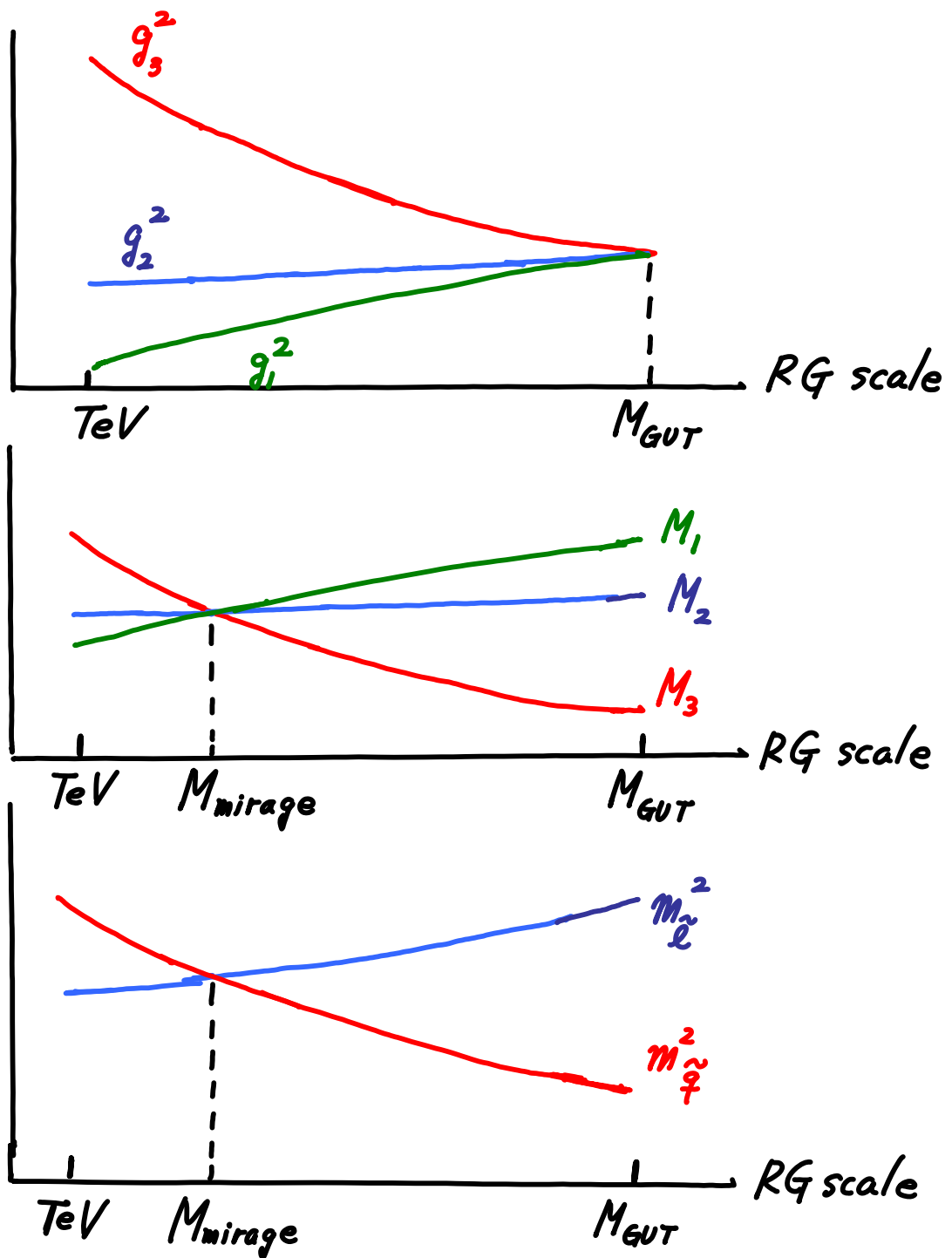
$$M_{\text{mirage}} = M_{\text{GUT}} \left(m_{3/2} / M_{\text{Pl}} \right)^{\alpha/2}$$

$\equiv$ Mirage Messenger Scale

$$\alpha \equiv \frac{m_{3/2}}{M_0 \underbrace{\ln(M_{\text{Pl}}/m_{3/2})}_{\sim 4\pi^2}} = 1 + \underbrace{l}_{\text{model-dependent rational number}}$$

$$\alpha = 1 \rightarrow M_{\text{mirage}} \sim \sqrt{m_{3/2} M_{\text{Pl}}} \quad (\text{intermediate mirage scale})$$

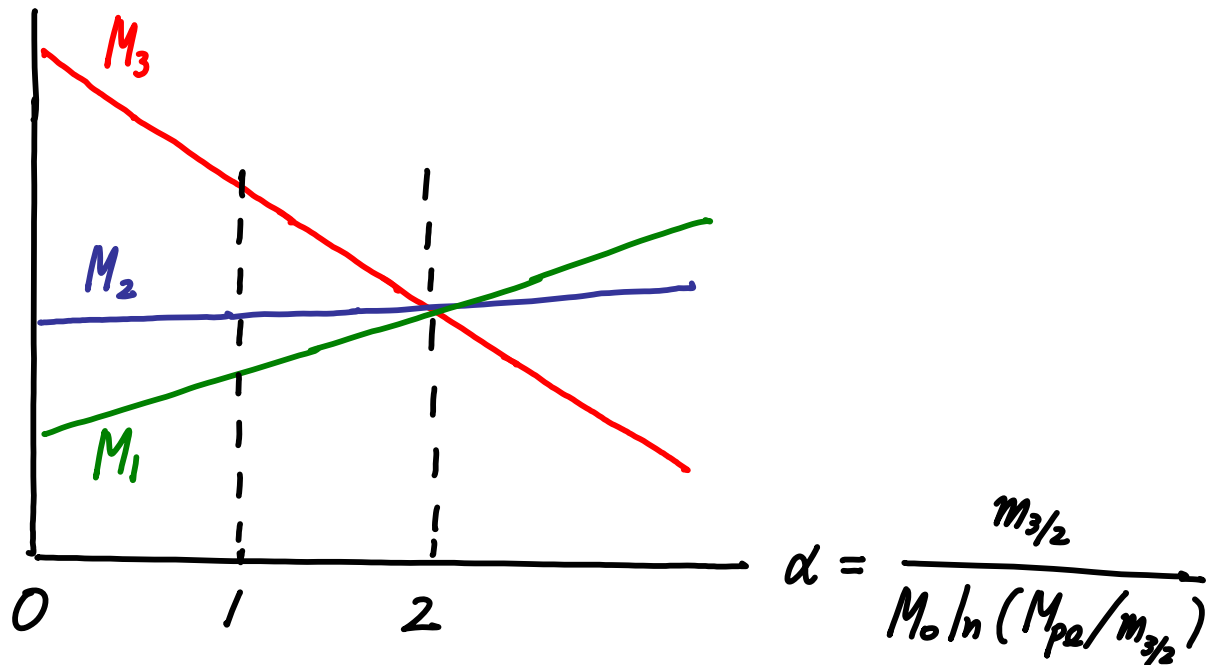
$$\alpha = 2 \rightarrow M_{\text{mirage}} \sim \text{TeV} \quad (\text{TeV mirage scale})$$



Compressed SUSY spectrum at TeV !

KC, Jeong, Okumura ; Endo, Yamaguchi, Yoshioka

Low energy gaugino masses as a ftn of α



$$M_1 : M_2 : M_3$$

$$1 : 2 : 6$$

$\alpha = 0$
(mSUGRA, GMSB)

$$1 : 1.3 : 2.4$$

$\alpha = 1$
(intermediate scale
mirage mediation)

$$1 : 1 : 1$$

$\alpha = 2$
(TeV scale mirage)

$$3.3 : 1 : 10$$

$\alpha = \infty$
(AMSB)

◆ μ & B , EWSB, ~~CP~~, FCNC, Cosmology

- μ & B (KC, Jeong, Okumura)

$$\alpha \neq 2 \rightarrow \text{NMSSM}$$

$$\alpha = 2 \rightarrow W_\mu = \kappa e^{-8\pi^2(kT - lS_0)} H_u H_d$$

- little SUSY fine-tuning

$$\Delta m_{H_u}^2 = -\frac{3y_t^2}{4\pi^2} m_{\tilde{t}}^2 \ln \left(\frac{M_{\text{mirage}}}{m_{\tilde{t}}} \right)$$

Mirage mediation generically ameliorates the little fine tuning for $m_h > 114 \text{ GeV}$.

$\alpha = 2$ ($M_{\text{mirage}} \sim \text{TeV}$) is the best scenario for the little fine tuning problem.
(KC, Jeong, Kobayashi, Okumura ; Kitano, Nomura)

- SUSY CP problem

Axionic shift symmetry

$$U(1)_T : \text{Im}(T) \rightarrow \text{Im}(T) + \text{real constant}$$

KC (1994)

assures that soft preserve CP :

$$K = K_0(T + \bar{T}) + Z_i(T + \bar{T}) Q^i Q^i$$

$$W = \omega_0 + A e^{-8\pi^2 T} + \kappa S^3 + \lambda S H_u H_d + \dots$$

$$f_a = \underset{\substack{\uparrow \\ \text{real constant}}}{\kappa_a T} + \text{constant}$$

- Using $U(1)_T$, $U(1)_R$ & other appropriate field redefinition, one can always make

$\omega_0, A, \kappa, \lambda$ real.

- $U(1)_T$ also assures $\partial_T K, \partial_T Z_i, \partial_T f_a, \dots$ are all real.

→ All soft parameters are real.

(KC, Falkowski, Nilles, Olechowski ; Endo, Yamaguchi, Yoshioka)

- SUSY flavor problem

Typically $Z_i = \frac{1}{(T + \bar{T})^{n_i}}$

(n_i = discrete modular weights)

If n_i are flavor universal, which is a rather *natural possibility*, soft terms preserve flavor also.

- Moduli & gravitino cosmology & DM

Still we might need a mechanism to eliminate coherent moduli oscillation.

(M. Yamaguchi's talk)

Variety of possibilities for LSP DM.

(K.I. Okumura's talk)

◆ Summary

- String compactification stabilizing moduli at phenomenologically viable dS vacuum with low energy SUSY suggests (or predict) *mirage mediation* with a *mirage messenger scale hierarchically lower than M_{pl}* which appears as a consequence of dynamical cancellation between moduli mediation & anomaly mediation.

- *Mirage mediation*

i) gives highly distinctive low energy superparticle spectrum :

Compressed spectrum which is clearly distinguishable from $mSUGRA$, $GMSB$ & $AMSB$.

ii) naturally avoid dangerous CP & $FCNC$ & also ameliorate (even avoid) the little fine-tuning for $EWSB$

iii) allows various types of LSP DM

(Next talk by K. Okumura)

